

# Applications of the Monte Carlo method

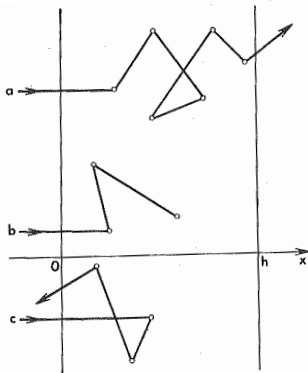
## Examples

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# Neutron transmission through a plate

The flux of neutrons with energy  $E_0$  is incident on a homogeneous infinite plate  $0 \leq x \leq h$ . The angle of incidence is  $90^\circ$ . Upon collision with atoms of the plate material neutrons may be either elastically scattered or captured. Possible fates of a neutron are depicted below: (a) it either passes, (b) is captured, or (c) is reflected by the plate.



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We want to calculate the probability of neutron transmission through the plate  $p^+$ , the probability of neutron reflection by the plate  $p^-$  and the probability of neutron capture inside the plate  $p^0$ .

Let assume for simplicity that energy of a neutron is not changed in scattering and that any direction of *recoil* of a neutron from an atom is equally probable, which is actually the case in neutron collisions with heavy atoms.

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Interaction of neutrons with matter is characterized by two constants: the capture cross section  $\sigma_c$  and the scattering cross section  $\sigma_s$ . The sum of them called the *total cross section*  $\sigma$

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The *free path length of a neutron*  $\lambda$ , i.e. the length of the path from one collision to another is a random variable with the probability density

$$p(x) = \sigma e^{-\sigma x}.$$

Let us check the normalization condition, assuming that the macroscopic width of the plate  $h$  can be set  $\infty$ .

$$\int_0^{\infty} p(x) dx = \sigma \int_0^{\infty} e^{-\sigma x} dx == \sigma \left[ -\frac{1}{\sigma} e^{-\sigma x} \right]_0^{\infty} = \sigma \frac{1}{\sigma} e^0 = 1,$$

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Similarly, let us find the expectation value of  $\lambda$

$$\begin{aligned} M\lambda &= \sigma \int_0^{\infty} x e^{-\sigma x} dx = \left\{ \begin{array}{l} u = x \Rightarrow du = dx \\ dv = e^{-\sigma x} dx \Rightarrow v = -\frac{1}{\sigma} e^{-\sigma x} \end{array} \right\} \\ &= \sigma \left[ -\frac{x}{\sigma} e^{-\sigma x} \Big|_0^{\infty} + \frac{1}{\sigma} \int_0^{\infty} e^{-\sigma x} dx \right] = -\frac{1}{\sigma} e^{-\sigma x} \Big|_0^{\infty} = \frac{1}{\sigma}. \end{aligned}$$



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Let us solve the formula for drawing  $\lambda$  values

$$\sigma \int_0^{\lambda} e^{-\sigma x} dx = \gamma,$$

where  $\gamma$  is the random variable uniformly distributed in the interval  $(0, 1)$ .

$$\sigma \left[ -\frac{1}{\sigma} e^{-\sigma x} \right]_0^{\lambda} = \gamma \Rightarrow -e^{-\sigma \lambda} + 1 = \gamma \Rightarrow e^{-\sigma \lambda} = 1 - \gamma.$$

Hence

$$\lambda = -\frac{1}{\sigma} \ln(1 - \gamma).$$

However,  $1 - \gamma$  is also random variable uniformly distributed in the interval  $(0, 1)$ , the same as  $\gamma$ , therefore we can write

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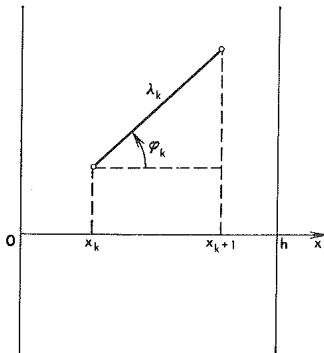
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# Calculation of a definite integral

We select a random direction of the neutron after scattering. As the problem is symmetric with respect to rotations about the  $Ox$ -axis, the neutron direction after  $k$ -th scattering inside the plate at the point with abscissa  $x_k$  is completely determined by the angle  $\varphi_k$ , as in the figure below.



We are now ready to simulate the trajectory of a neutron.

# Neutron transmission through a plate

Actually, we can use the random variable  $\mu_k = \cos \varphi_k$  instead, which we assume to be uniformly distributed in the interval  $(-1, 1)$ . Thus, we generate

$$\mu_k = \cos \varphi_k = [1 - (-1)]\gamma - 1 = 2\gamma - 1.$$

We draw the free path length

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Then we check the condition of neutron transmission through the plate

$$x_{k+1} > h.$$

If it is so, the computation of the trajectory is terminated and 1 is added to the counter of transmitted neutrons.

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# Neutron transmission through a plate

Otherwise we check the condition of reflection

$$x_{k+1} < 0.$$

If it is so, the computation of the trajectory is terminated and 1 is added to the counter of reflected neutrons.

If neither of the two conditions are satisfied, i.e.

$$0 \leq x_{k+1} \leq h,$$

which means that the neutron has undergone the  $(k + 1)$ -th collision inside the plate, we choose another value of  $\gamma$  and check the capture condition

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Obviously, the initial values for each trajectory are

$$x_0 = 0, \quad \cos \varphi_0 = 1.$$

After  $N$  trajectories are sampled we obtain the following results

- $N^+$  neutrons were transmitted through the plate,
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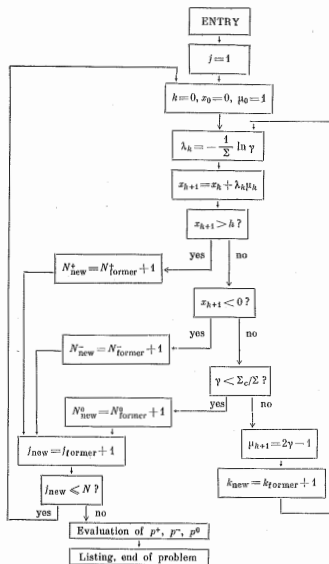
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Here is the block scheme of the computer program for this problem





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Let us consider a function  $g(x)$  defined in the interval  $a < x < b$ .  
We want to approximate the integral

$$I = \int_a^b g(x) dx.$$

Although this problem is not at all probabilistic, we will apply the Monte Carlo method to solve it.

Let us choose an arbitrary distribution density  $p_\xi(x)$  specified in the interval  $(a, b)$ , i.e. an arbitrary function  $p_\xi(x)$  satisfying the conditions

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Along with the random variable  $\xi$  defined on the interval  $(a, b)$  with the density  $p_\xi(x)$ , we will need another random variable

$$\eta = \frac{g(\xi)}{p_\xi(\xi)}.$$

Let us calculate the expectation value of  $\eta$

$$M\eta = \int_a^b \left[ \frac{g(x)}{p_\xi(x)} \right] p_\xi(x) dx = I.$$

Let us consider now  $N$  identical independent random variables  $\eta_1, \eta_2, \dots, \eta_N$  and apply the central limit theorem to their sum. Then we will obtain the relation

$$P \left\{ \left| \frac{1}{N} \sum_{j=1}^N \eta_j - I \right| < 3 \sqrt{\frac{D\eta}{N}} \right\} \approx 0.997.$$

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This relation means that if we sample  $N$  values  $\xi_1, \xi_2, \dots, \xi_N$ , then for sufficiently large  $N$

$$I \approx \frac{1}{N} \sum_{j=1}^N \frac{g(\xi_j)}{p_{\xi}(\xi_j)}.$$

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Independent of which random variable  $\xi$ , defined in the interval  $(a, b)$ , we use, we will obtain

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# Calculation of a definite integral

However, the variance  $\mathbf{D}\eta$ , and hence the estimated error of our approximation, depends on what specific variable  $\xi$  is used. Indeed

$$\mathbf{D}\eta = \mathbf{M}\eta^2 - I^2 = \int_a^b \left[ \frac{g(x)}{p_\xi(x)} \right]^2 p_\xi(x) dx - I^2 = \int_a^b \frac{g^2(x)}{p_\xi(x)} dx - I^2.$$

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# The choice of the best $\xi$

To find the minimum of the variance

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among all possible choices of  $p_\xi(x)$  we will use the Schwartz inequality which for real functions  $u(x)$  and  $v(x)$  integrable with their squared module in the interval  $(a, b)$  has the form

$$\left[ \int_a^b |u(x)v(x)| dx \right]^2 \leq \int_a^b u^2(x) dx \int_a^b v^2(x) dx.$$

Let us set  $u(x) = \frac{g(x)}{\sqrt{p_\xi(x)}}$  and  $v(x) = \sqrt{p_\xi(x)}$ , then we will obtain

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